



SIMULATING EXTREME SPACETIMES

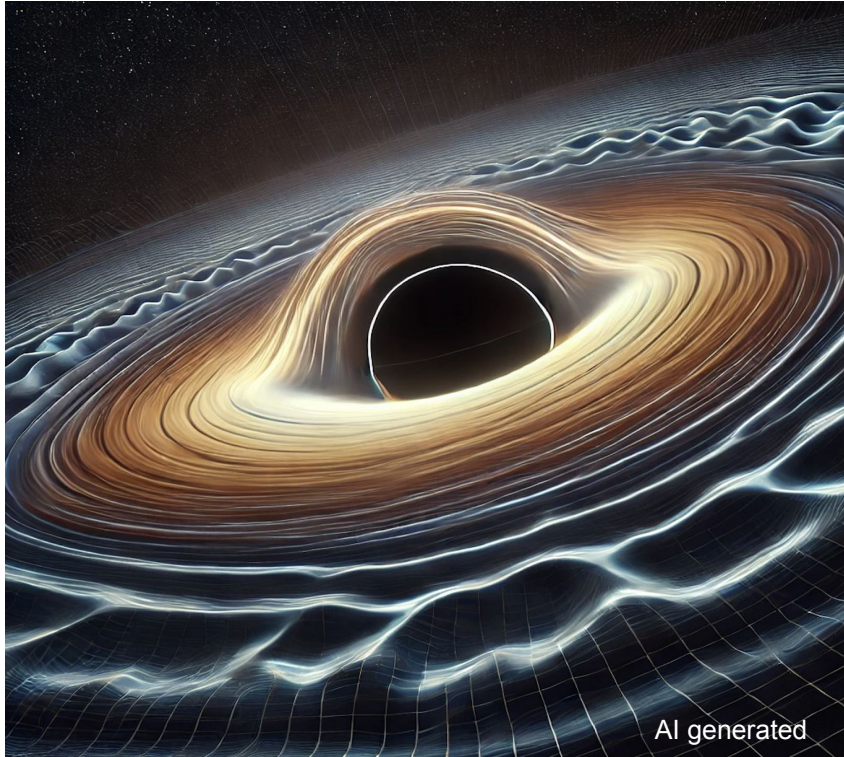
Black holes, neutron stars, and beyond...

Caltech

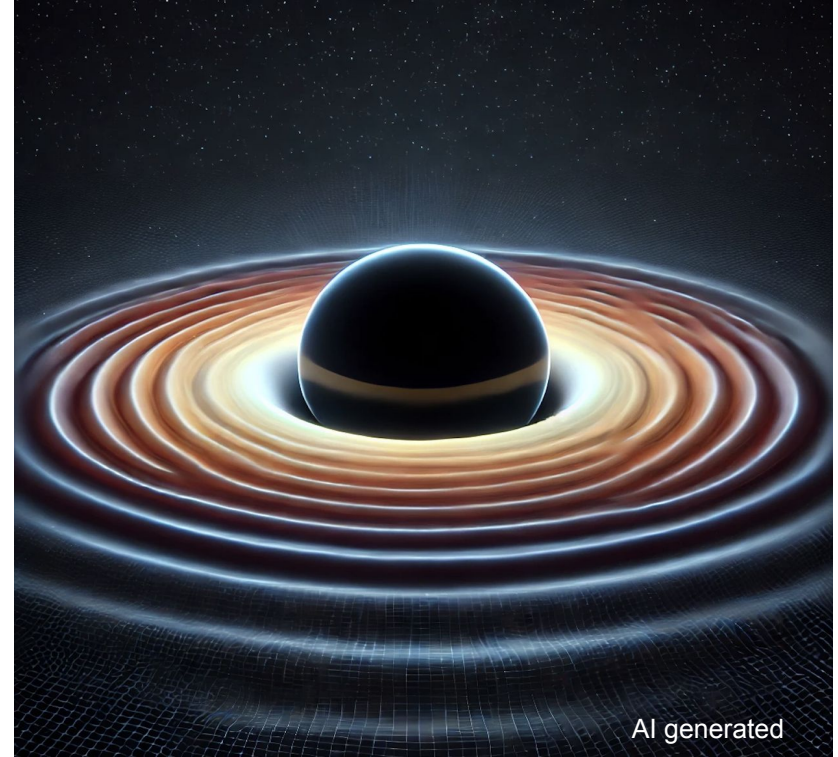
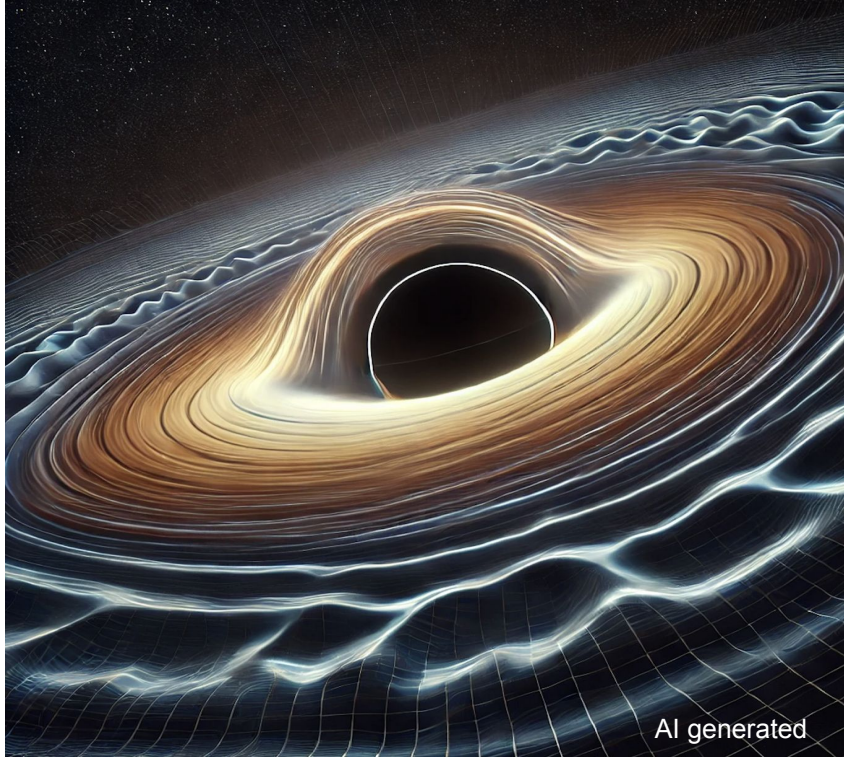
RINGDOWN IN IN BINARY BLACK HOLE SIMULATIONS

Isabella Pretto

Ringdown



Ringdown



Ringdown

- 1957: Birth of BH perturbation theory (Regge and Wheeler)
- 1963: Kerr solution is discovered
- 1967: Wheeler coins the term 'black hole'
- 1970-71: Schwarzschild BHs (Regge, Wheeler, Zerilli, Vishveshwara, Press, Davis...)
- 1973: Teukolsky separates the eqs. in Kerr
- 1975: Characteristic frequencies (Chandrasekhar, Detweiler...)

Ringdown - Kerr

- Kerr metric: most general vacuum BH sol. in 4-D asymp. flat
- Curvature invariants $\rightarrow$ perturbation equations decouple and separate
- Master perturbation eq. for fields of general spin
- Newman-Penrose formalism

$$\Psi_0 = -C_{1313} = -C_{\mu\nu\lambda\sigma} l^\mu m^\nu l^\lambda m^\sigma,$$

$$\Psi_4 = -C_{2424} = -C_{\mu\nu\lambda\sigma} n^\mu m^{*\nu} n^\lambda m^{*\sigma}.$$

Ringdown - Kerr

➤ For a spin- s field,

$$\psi(t, r, \theta, \phi) = \frac{1}{2\pi} \int e^{-i\omega t} \sum_{l=|s|}^{\infty} \sum_{m=-l}^l e^{im\phi} {}_sS_{lm}(\theta) R_{lm}(r) d\omega,$$

Ringdown - Kerr

$$\left[\frac{\partial}{\partial u} (1 - u^2) \frac{\partial}{\partial u} \right]_s S_{lm} + \left[a^2 \omega^2 u^2 - 2a\omega s u + s +_s A_{lm} - \frac{(m + su)^2}{1 - u^2} \right]_s S_{lm} = 0,$$

$$\Delta \partial_r^2 R_{lm} + (s + 1)(2r - 2M) \partial_r R_{lm} + V R_{lm} = 0.$$

$$V = 2is\omega r - a^2\omega^2 -_s A_{lm} + \frac{1}{\Delta} [(r^2 + a^2)^2 \omega^2 - 4Mam\omega r + a^2 m^2 \\ + is(am(2r - 2M) - 2M\omega(r^2 - a^2))].$$

Ringdown - Kerr

- Quasinormal modes (QNMs) and boundary conditions

$$\Psi \sim e^{-i\omega(t+r_*)}, \quad r_* \rightarrow -\infty (r \rightarrow r_+).$$

$$\Psi \sim e^{-i\omega(t-r_*)}, \quad r \rightarrow \infty,$$

Ringdown

- After prompt response, quasinormal modes (QNMs)

$$h(u, \theta, \phi) = \sum_j \left\{ \begin{matrix} \ell \geq 2, |m| \leq \ell, \\ n \geq 0, p = \pm 1 \end{matrix} \right\} C_j e^{-i\omega_j t} {}_{-s}S_{(\ell, m)}(a\omega_j)$$

- $\omega_{(l, m, n, p)}$ depends only on mass and spin of final black hole
- Detect at least 2 frequencies— test the no-hair theorem

Ringdown

- Second-order perturbation

$$\mathcal{T}\psi = \mathcal{S},$$

$$\omega \equiv \omega_{(\ell_1, m_1, n_1)} + \omega_{(\ell_2, m_2, n_2)}.$$

Fitting procedure

$$h(u, \theta, \phi) \sim \sum_j = \left\{ \begin{array}{l} \ell \geq 2, |m| \leq \ell, \\ n \geq 0, p = \pm 1 \end{array} \right\} C_j e^{-i\omega_j t}$$

Fitting procedure

- Well-posedness
 - The solution exists
 - It is unique
 - It depends uniquely on the data
- Solving for QNMs is an ill-conditioned problem
- Variable projection algorithm

$$f(\omega, C, t) = c_1\phi(\omega_1, t) + c_2\phi(\omega_2, t) + \dots + c_n\phi(\omega_n, t)$$

Stability

➤ **Measured** QNM amplitude should decay as $\sim e^{t \times \text{Im}[\omega]}$

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- Fractional variation of C_{QNM} at reference time

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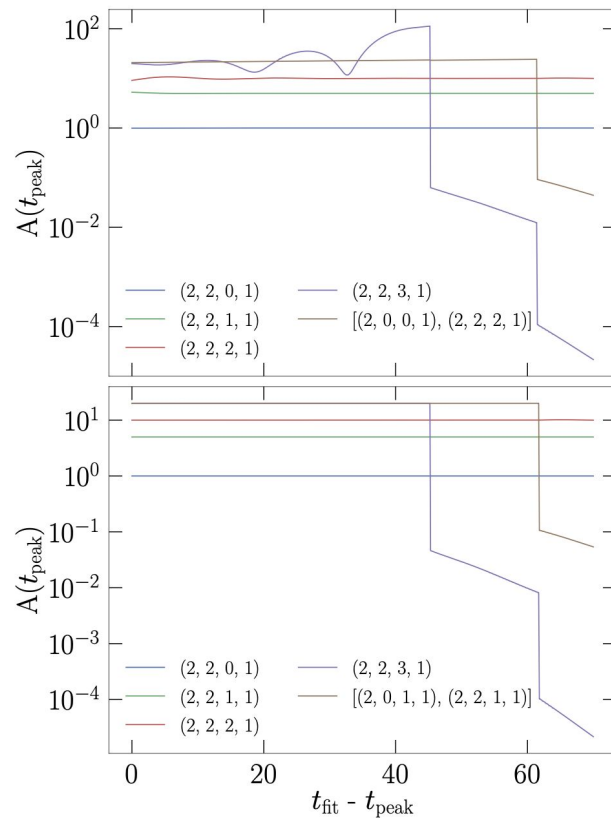
$$\sigma_{\text{QNM}} \stackrel{\text{def}}{=} \sqrt{(\Delta \text{Re}[C_{\text{QNM}}]/|C_{\text{QNM}}|)^2 + (\Delta \text{Im}[C_{\text{QNM}}]/|C_{\text{QNM}}|)^2}$$

- Require $\sigma_{\text{QNM}} < \sigma_{\text{max}}$ over $\Delta t_{\text{QNM}} = -\ln(\Delta A_{\text{min}})/\text{Im}[\omega_{\text{QNM}}]$

Stability

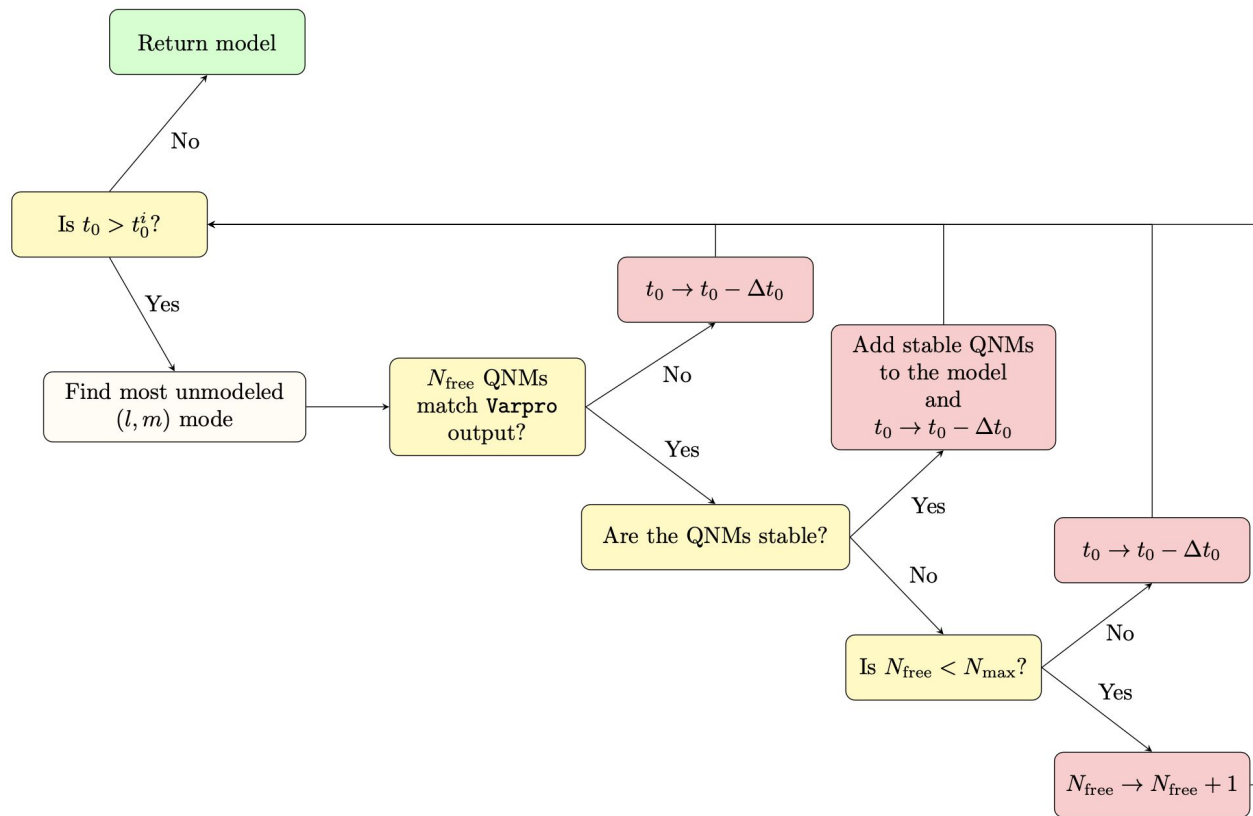
$(2,0,1,1) \times (2,2,1,1)$: $0.807 - 0.528i$

$(2,0,0,1) \times (2,2,2,1)$: $0.801 - 0.530i$

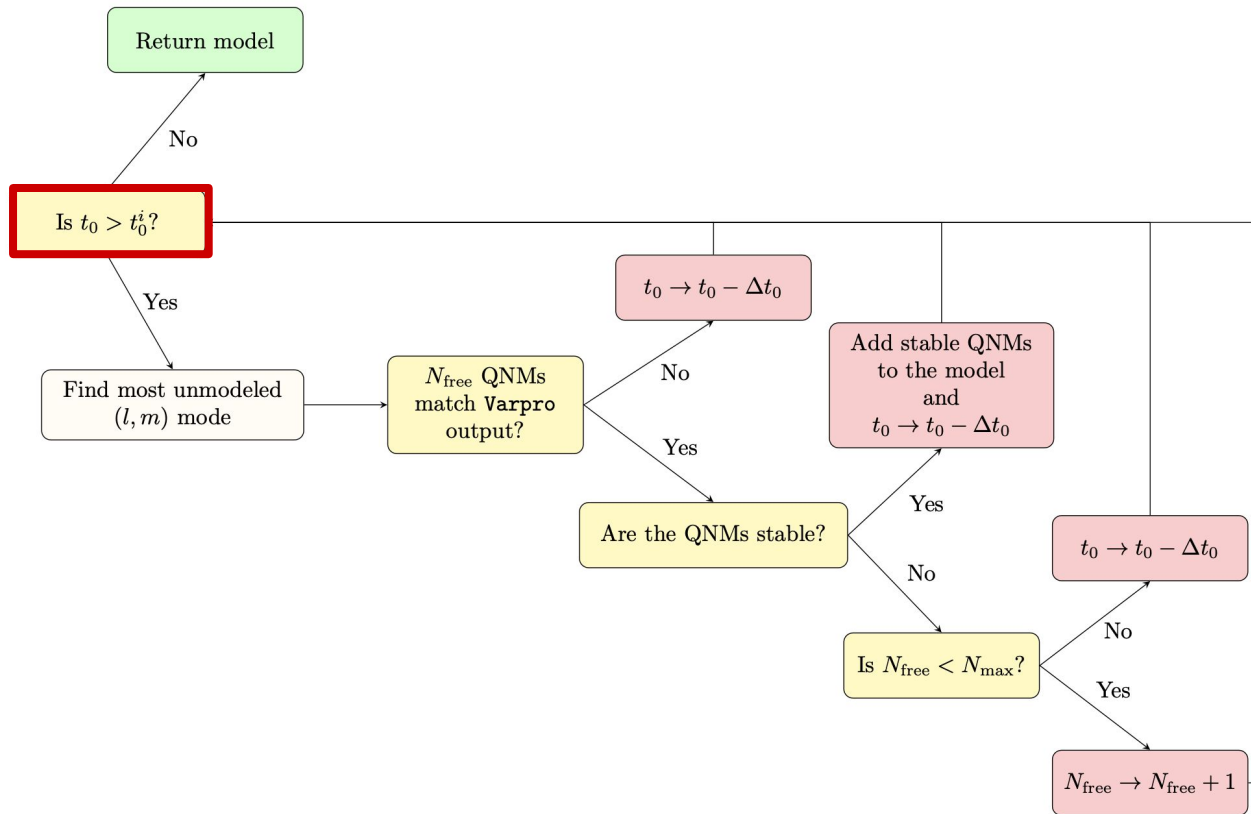


How to automate the search for QNMs in numerical relativity waveforms?

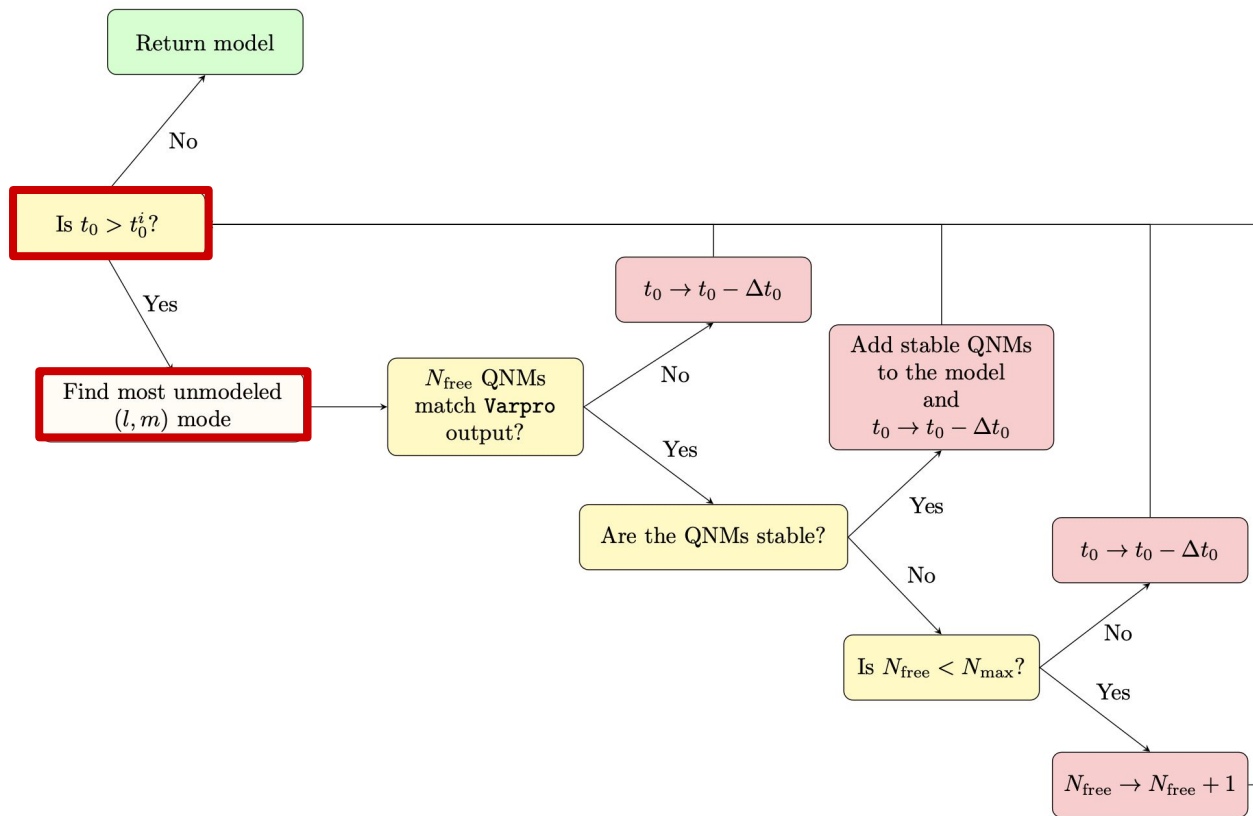
Building a QNM model



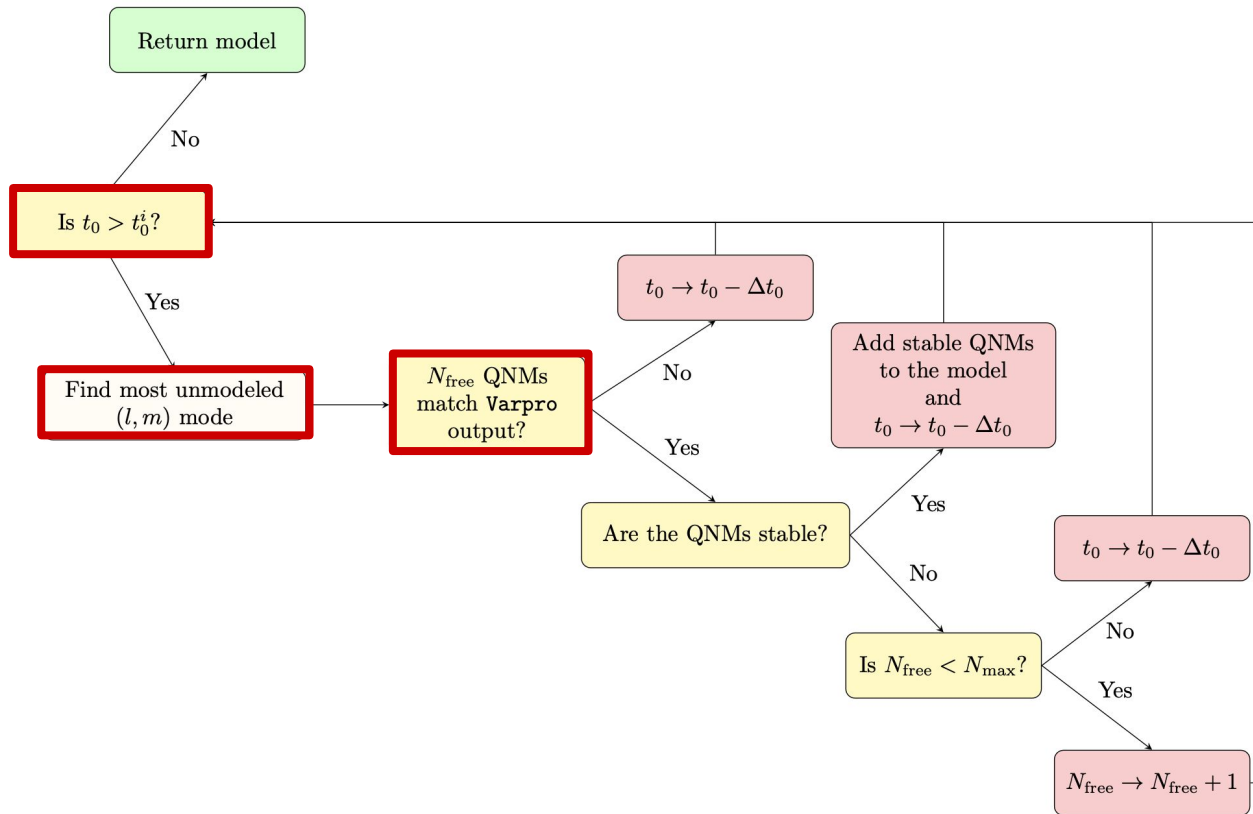
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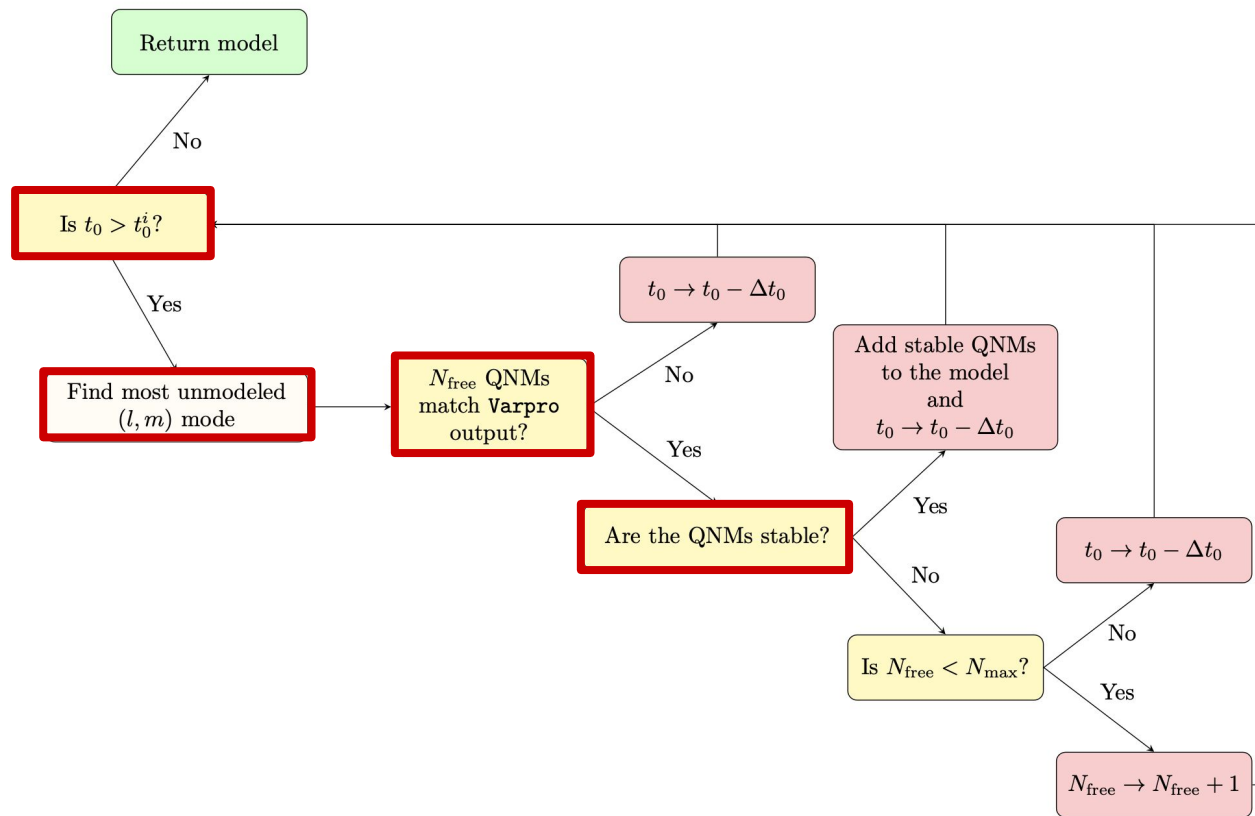
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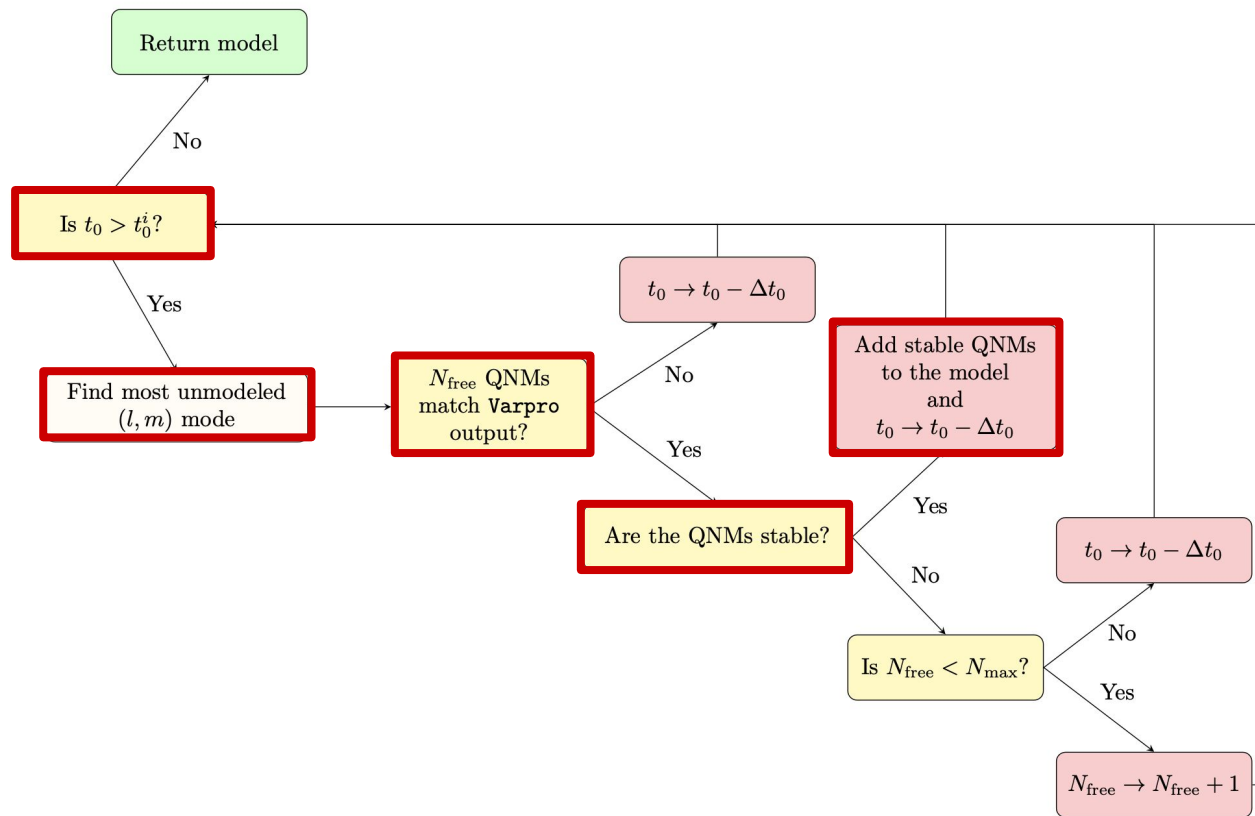
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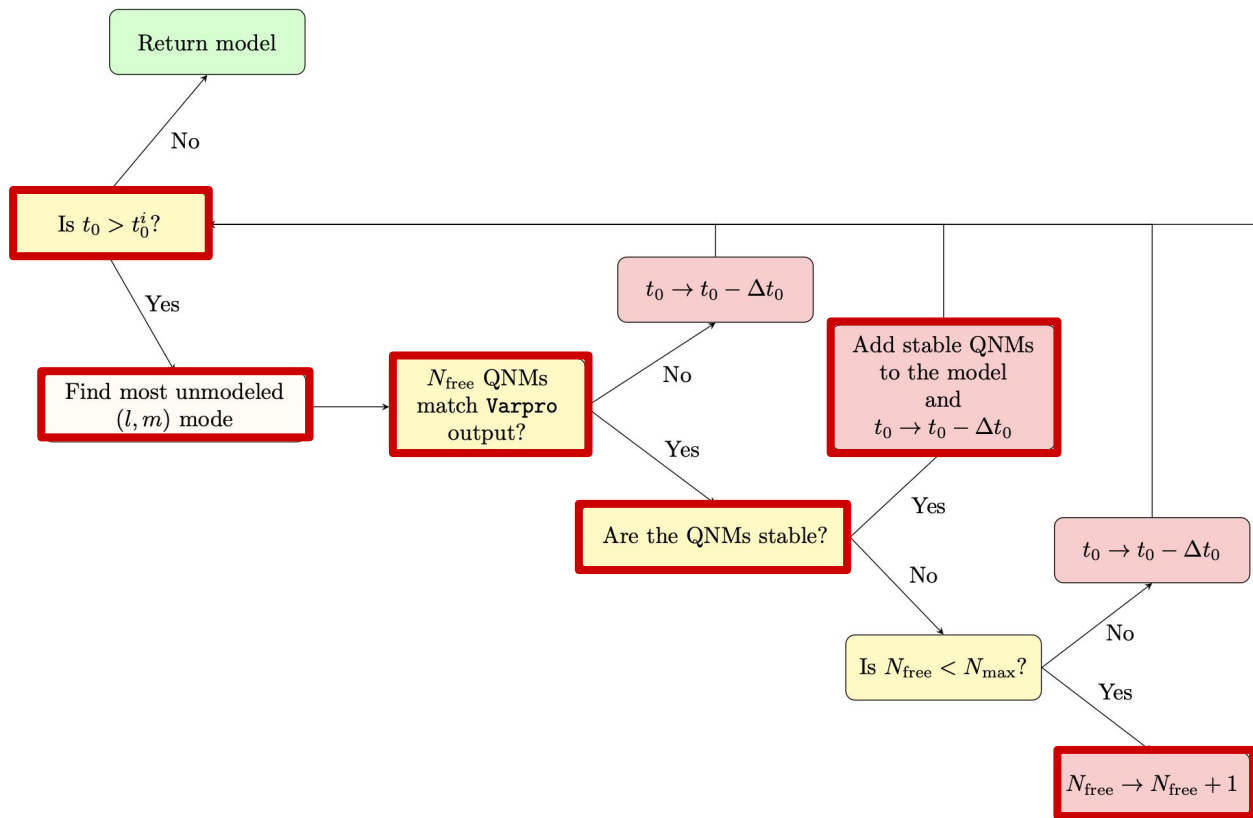
Building a QNM model



Building a QNM model



Building a QNM model



Features

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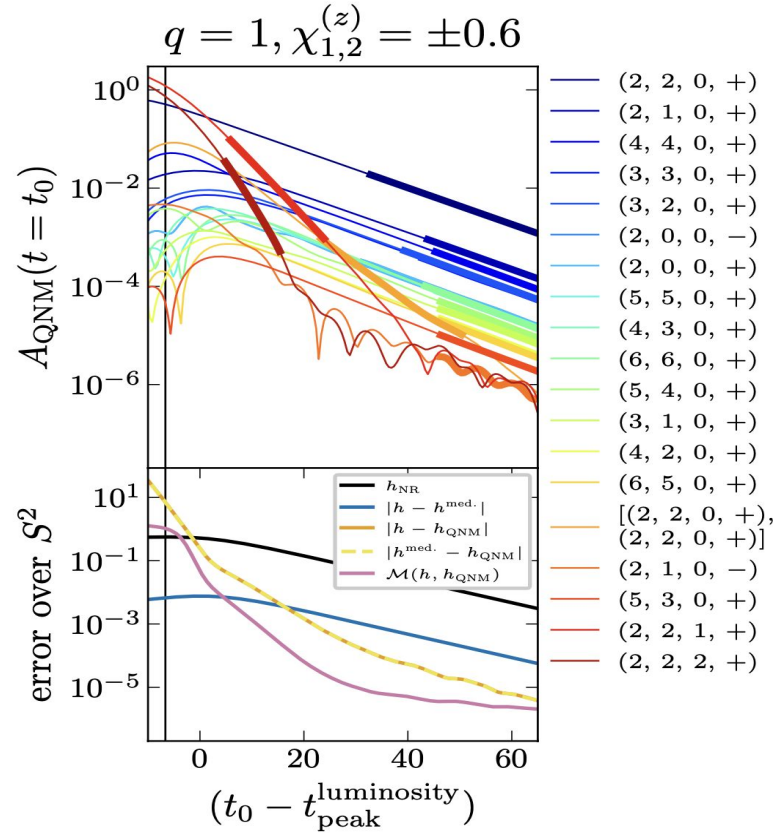
Tests

- > 55 mock waveforms, with non-Gaussian/Gaussian noise

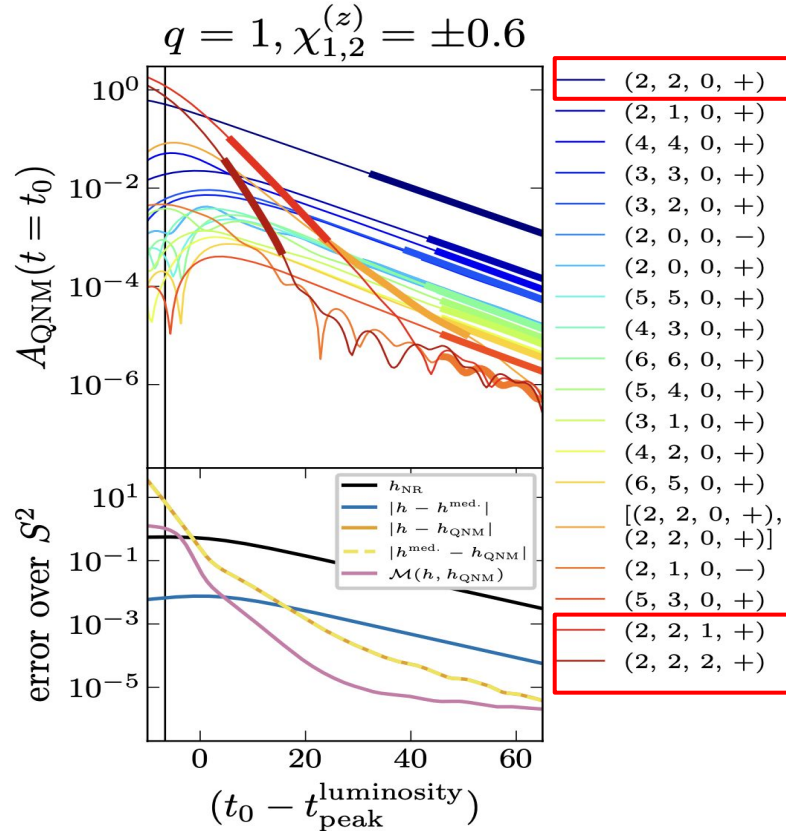
$$\sigma_{\text{noise}} \in \{10^{-6}, 10^{-5}, \dots, 0.1\}$$

- Find optimal values for A_{min} and σ_{max}
- Validate algorithm

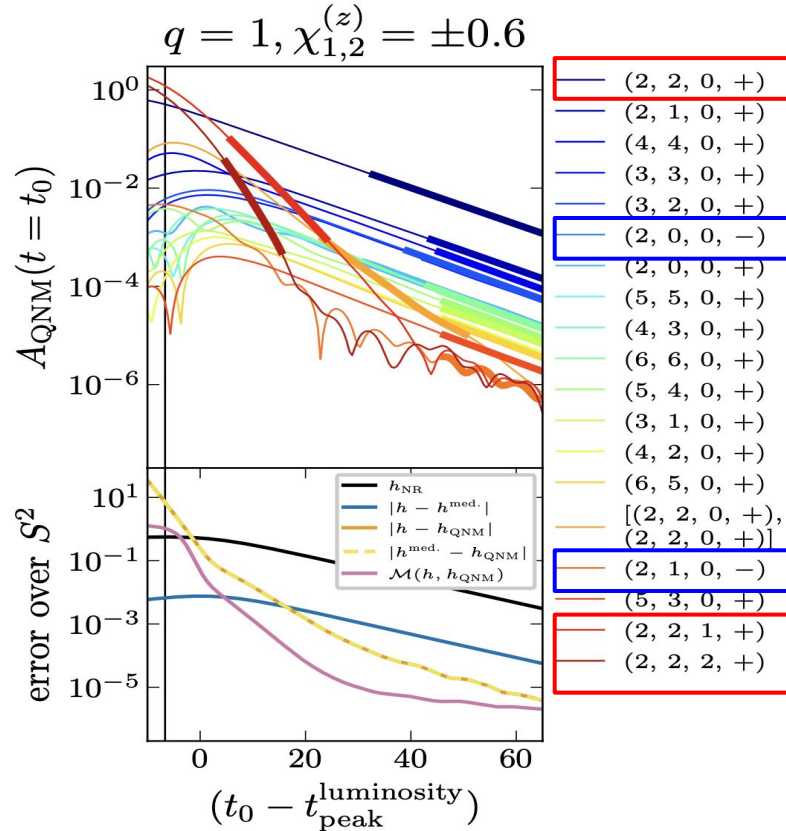
NR waveforms



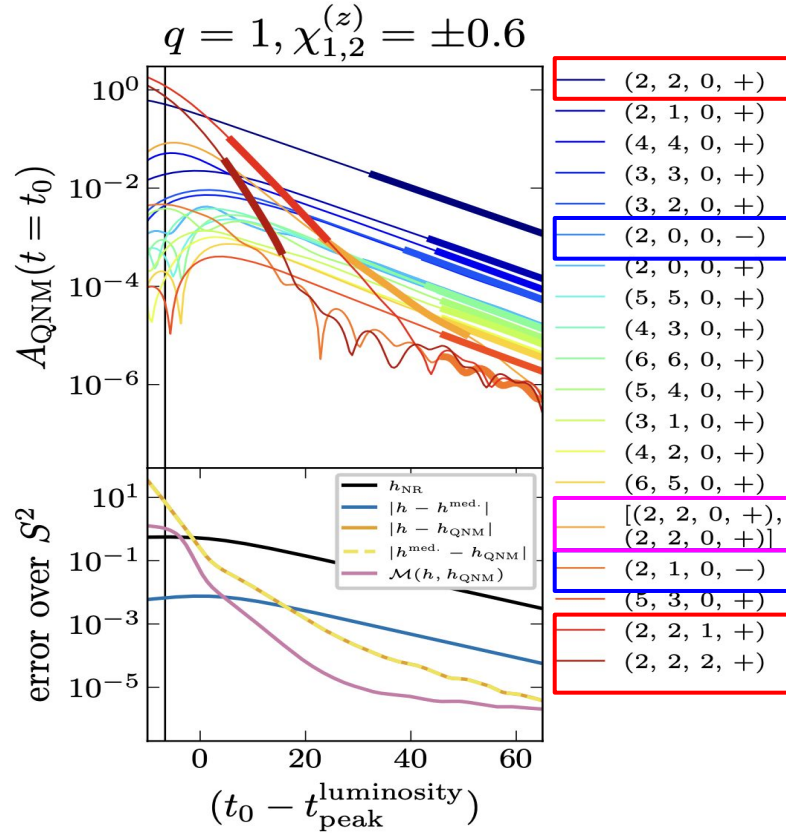
NR waveforms



NR waveforms



NR waveforms



NR waveforms

- Analyze the excitation of QNMs
- ~400 quasi-circular, non-precessing systems with $q < 8$, $|\chi^{(z)}_{1,2}| < 0.8$

- $$C_{\text{QNM}} = I_{\text{QNM}} E_{\text{QNM}},$$



Source
function



Excitation
factor

NR waveforms

- Analyze the excitation of QNMs
- ~400 quasi-circular, non-precessing systems with $q < 8$, $|\chi_{1,2}^{(z)}| < 0.8$

➤ ~~$C_{\text{QNM}} = I_{\text{QNM}} E_{\text{QNM}}$~~ $E_{\text{QNM}} = \frac{1}{I_{(\ell,m)}} e^{-i\omega_{\text{QNM}} \Delta t_0^{(\ell,m)}} C_{\text{QNM}}$

C_{QNM}  Source function

I_{QNM}  Excitation factor

Oshita, Cardoso (2024)
arXiv:240702563

NR waveforms

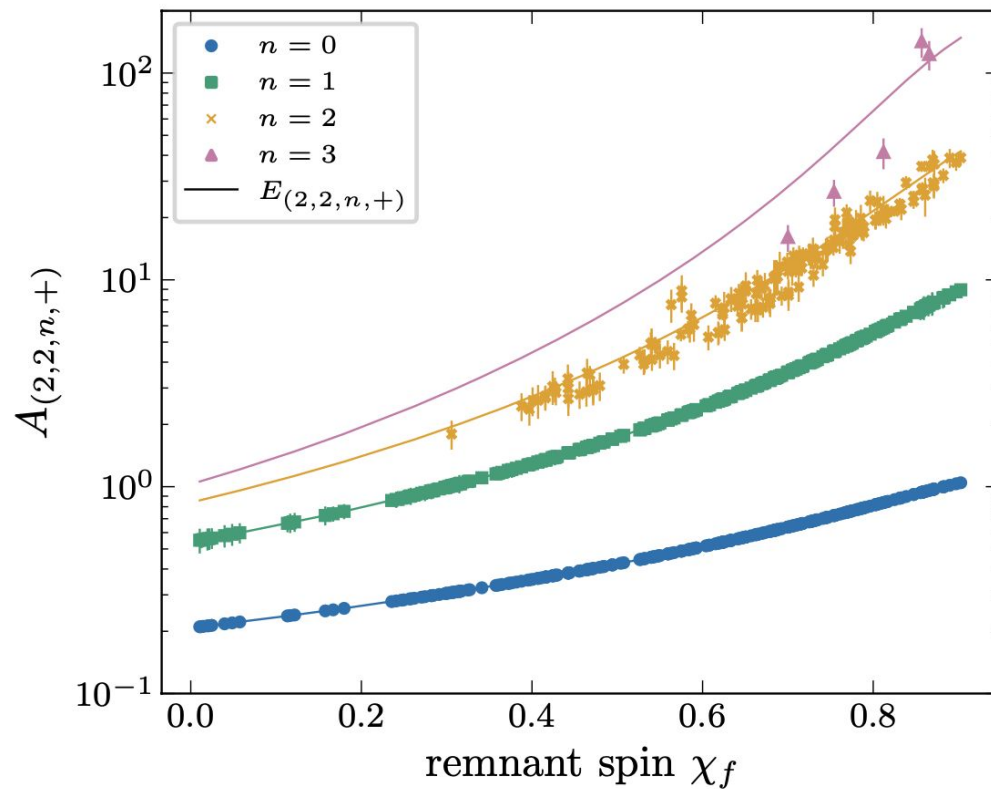
- Source function— flat for $(2, 2, n, +)$ overtones?
- Constrain free parameters

Oshita (2021)
arXiv:2109.09757

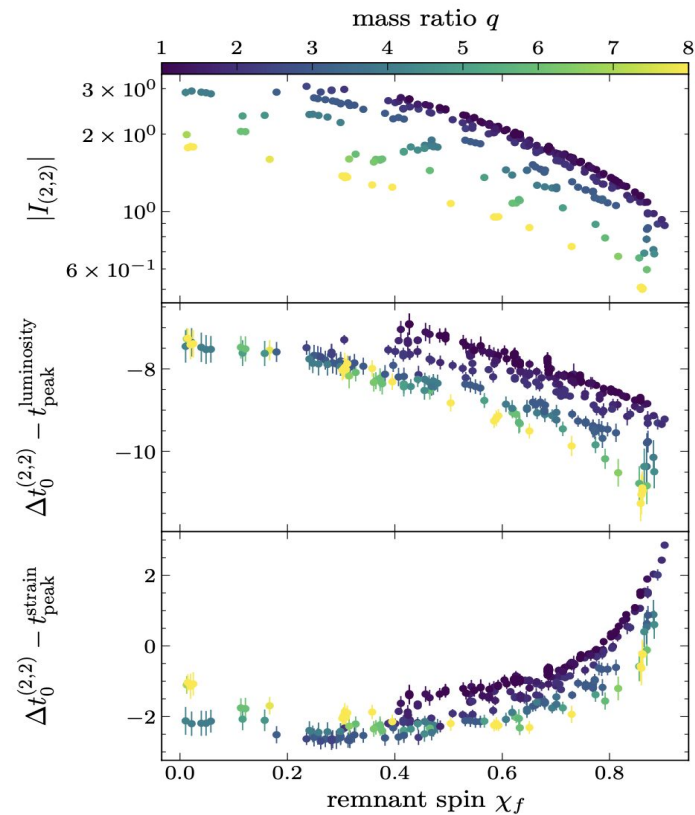
Giesler et al. (2024)
arXiv:2411.11269

$$\Delta t_0^{(2,2)} = \frac{\ln \left(\left| \frac{E_{(2,2,0,+)} C_{(2,2,1,+)}}{E_{(2,2,1,+)} C_{(2,2,0,+)}} \right| \right)}{\text{Im}[\omega_{(2,2,0,+)} - \omega_{(2,2,1,+)}]}$$
$$I_{(2,2)} = \frac{C_{(2,2,0,+)}}{E_{(2,2,0,+)}} e^{-i\omega_{(2,2,0,+)} \Delta t_0}$$

NR waveforms



NR waveforms



Conclusion

- Iterative fitting with Varpro + on-the-fly stability analysis
- Mock and NR waveforms
- Physics of asymmetric systems and (2,2) excitation
- Future 
 - Different BBH systems
 - Breakdown of constant-QNM model?